On the wave-current interactions



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# Some informations about the wave-current coupling...

#### 1. Coupled wave-current modelling

# a. The 3D challenge

Momentum equation formulated for:

- total momentum (includes Stokes drift): this is too complex (vertical flux of wave momentum is a strange beast)
- mean flow momentum only





#### 1. Wave-current 3D coupling

#### **b.** Summary

#### Where we are now:

# > 2-way coupling of WAVEWATCH III and MARS3D with PALM

- \* Based on WWATCH version 3.14\_Ifremer
  & SHOM and MARS3D version 8.0
  \* Wrong equations (Mellor 2003 and 2008) well implemented
- \* Correct equations implemented realistic validation OK (beach)

# > Now working on:

- Mixing/friction parameterizations
- Rip currents





Impact of the mixing due to the wave breaking on the bottom friction ... and the consequences

a. Objectives

\* To evaluate the impact of wave breaking on the bottom friction and the consequences on the longshore current and on the set-up.

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 $\Rightarrow$  The parametrization of Mellor (2002) is used: the bottom friction depends on the turbulent kinetic energy.

i) To redo the numerical experiments presented in Mellor (2002) paper
 + Addition of wave breaking at the surface.

=> comparison between the phase-averaged case and the phaseresolving case

- ii) Application in nearshore zone:
  - ML02 vs Walstra (2000)+Soulsby (1995).
  - Impact on the longshore current
  - Impact on the set-up



# **b.** Phase-averaged vs Phase-resolving

Equations for the phase-resolving case

$$\frac{\partial u}{\partial t} = \frac{\tau_{0x}}{h} + u_{bx}\omega cos(\omega t) + \frac{1}{D}\frac{\partial \tau_x}{\partial \varsigma}, + \text{Soulsby (1995)}$$

$$\frac{\partial k}{\partial t} = \frac{1}{D^2} \cdot \frac{\partial}{\partial \varsigma} \left(\frac{\nu_V}{s_k} \cdot \frac{\partial k}{\partial \varsigma}\right) - \frac{\partial k}{\partial \varsigma} \cdot \frac{\partial \varsigma}{\partial t} + \text{Prod} + \text{Buoy} - \epsilon,$$

$$\frac{\partial \epsilon}{\partial t} = \frac{1}{D^2} \cdot \frac{\partial}{\partial \varsigma} \left(\frac{\nu_V}{s_\epsilon} \cdot \frac{\partial \epsilon}{\partial \varsigma}\right) - \frac{\partial \epsilon}{\partial \varsigma} \cdot \frac{\partial \varsigma}{\partial t} + \frac{\epsilon}{k} (c_1 \text{Prod} + c_3 \text{Buoy} - c_2 \epsilon F_{wall}).$$

Equations for the phase-averaged case

$$\begin{aligned} \frac{\partial \overline{u}}{\partial t} &= \boxed{\overline{\tau}_{0x}}{h} + \frac{1}{D} \frac{\partial \overline{\tau}_{x}}{\partial \varsigma}, \\ \frac{\partial \overline{k}}{\partial t} &= \frac{1}{D^{2}} \cdot \frac{\partial}{\partial \varsigma} \left( \frac{\overline{\nu}_{V}}{s_{k}} \cdot \frac{\partial \overline{k}}{\partial \varsigma} \right) - \frac{\partial \overline{k}}{\partial \varsigma} \cdot \frac{\partial \varsigma}{\partial t} + \overline{\mathrm{Prod}} + \overline{\mathrm{Buoy}} - \overline{\epsilon} + \overline{P_{k}}, \\ \frac{\partial \overline{\epsilon}}{\partial t} &= \frac{1}{D^{2}} \cdot \frac{\partial}{\partial \varsigma} \left( \frac{\overline{\nu}_{V}}{s_{\epsilon}} \cdot \frac{\partial \overline{\epsilon}}{\partial \varsigma} \right) - \frac{\partial \overline{\epsilon}}{\partial \varsigma} \cdot \frac{\partial \varsigma}{\partial t} + \frac{\overline{\epsilon}}{\overline{k}} \left( c_{1} \overline{\mathrm{Prod}} + c_{3} \overline{\mathrm{Buoy}} - c_{2} \overline{\epsilon} F_{wall} \right) + \overline{P_{eps}}. \\ \overline{P}_{eps} &= C \frac{\overline{\epsilon}}{\overline{k}} P_{k}. \qquad \overline{\tau}_{x}|_{z=0} = \frac{\overline{u} \kappa S_{M0} \sqrt{2\overline{k_{0}}}}{\ln \left(\frac{z}{z_{0}}\right)}. \qquad \overline{\tau}_{x}|_{z=0} = \frac{\overline{u} \kappa S_{M0} \sqrt{2\overline{k_{0}}}}{\ln \left(\frac{z}{z_{0}} + 1\right)}. \end{aligned}$$

### **b.** Phase-averaged vs Phase-resolving

- Oscillations of the wave bottom boundary layer with the wave phase for pure oscillatory flow. Vertical profiles of current



# **b.** Phase-averaged vs Phase-resolving

Mean flow superimposed on an oscillatory flow:

- \* Five meshes are tested:
- refined meshes with 1200 grid points.
- depth of the first grid point:
  - \*\* Mesh 1:  $z_{bot} = 3.0.10^{-2}$  m.
  - \*\* Mesh 2:  $z_{bot} = 9.2.10^{-4}$  m.
  - \*\* Mesh 3:  $z_{bot} = 3.2.10^{-8}$  m.
  - \*\* Mesh 4: z<sub>bot</sub> = 1.3.10<sup>-5</sup> m.
  - \*\* Mesh 5:  $z_{bot} = 7.6.10^{-5}$  m.

=> 0.2 < Fz(bottom) < 5.



# b. Phase-averaged vs Phase-resolving



• The original Fz function must be changed.

- Currents are very close.
- TKE near the bottom is greatly enhanced.



# c. Application in surf zone

\* NSTS configuration: Leadbetter beach

- \* Simulations of 3D circulation in surfzone: MARS-WWATCHIII
  - dx=4m, dy=20m
  - dt=1s
  - 100 sigma levels (refined mesh)

\* Hs<1.1m, T=12s, θ=109 deg

\* Impact of the modelization of the bottom shear stress on the longshore current and set-up

- Mellor (2002) vs Walstra (2000)+Soulsby(1995)



### c. Application in surf zone

The Fz function must be changed: => Inappropriated positive values near the surface





#### Hs is correctly simulated



# The ML02 parametrization must be modified to correctly simulate the current



# c. Application in surf zone

$$\tau_y|_{z=0} = A \cdot \widehat{v},$$

 $A = \frac{\kappa S_{M0}\sqrt{2k_0}}{\ln\left(\frac{z}{z_0}\right)}.$ \* TKE overestimated:

=> weakest longshore current

\* Modified bottom stress gives similar A term than in the WSB95 case:
⇒ longshore current is ok





3. Rip currents

# a. Configuration & Objectives

\* The bathymetry is an approximation of the beach profile measured at Duck on October 11, 1990

- \* Simulations of 3D circulation in surfzone: MARS-WWATCHIII
  - dx=12m, dy=12m, dt=1s
  - 15 sigma levels
- \* Hs=1m, T=10s, θ=90 deg
- \* Comparison of results between the one-way mode and the twomode
- barotropic currents
- vorticity
- forcing terms

# <u>One-way mode</u>: vector current over bathymetry



#### 3. Rip currents

# b. One-way mode vs Two-way mode: barotropic currents





#### 3. Rip currents

### d. Summary

# TO CONCLUDE:

\* Rip currents are correctly simulated by the 3D wave-current model

\* One way-mode vs Two-way mode

- The current intensity is reduced when the feedback is activated

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- The vorticity is reduced when the feedback is activated and its offshore extension is lesser

- The forcing terms from breaking are less important for the two-way mode than for the one-way mode

=> Here the impact of the feedback is weak

# **IN THE FUTURE:**

- \* High resolution numerical simulations
- \* To extend this study for other rip systems